



## Letter

Phase transition-induced high electromechanical activity in  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  lead-free ceramic system

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## ARTICLE INFO

## Article history:

Received 7 November 2008

Received in revised form 20 January 2009

Accepted 22 January 2009

Available online 6 February 2009

## Keywords:

Electromechanical activity

Lead-free ceramics

Morphotropic phase boundary (MPB)

Phase transition

## ABSTRACT

A high electromechanical activity is observed in the  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) lead-free ceramic system at and around the orthorhombic (*O*)-tetragonal (*T*) phase transition temperature ( $T_{O-T}$ ). This activity is found to originate from an *O-T* phase transition region at ambient temperature rather than a classical morphotropic phase boundary (MPB) region intrinsic in the  $Pb(Zr_{1-x}Ti_x)O_3$  (PZT) lead-based ceramic system. Li modification enables a large decrease in  $T_{O-T}$  instead of constituting a classical MPB. In contrast to the nearly temperature-independent classical MPB behavior in the PZT system, the strong temperature-dependent phase transition behavior in the system may impose a challenge to temperature demanding applications.

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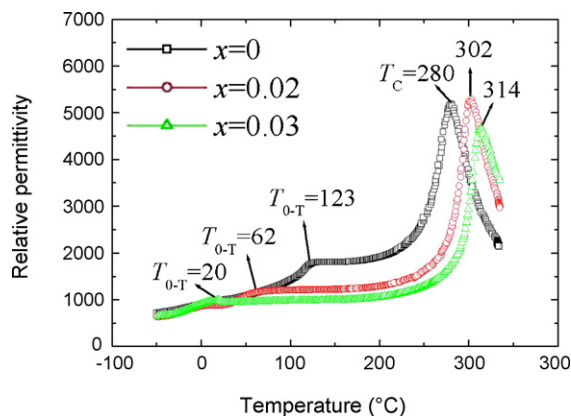
Traditional piezoelectric ceramic systems, such as  $Pb(Zr_{1-x}Ti_x)O_3$  (PZT) and  $(1-x)Pb(Mg_{1/3}Nb_{2/3})O_3-xPbTiO_3$  (PMN-PT), due to the constitution of large amounts of toxic lead (Pb), will be totally restricted from “green” products in the near future from the viewpoint of environmental protection [1,2]. Accordingly, the development of environmentally friendly lead-free ceramic systems has been an urgent demand in the recent decade. Among all the possible lead-free systems, the  $(K_{0.5}Na_{0.5})NbO_3$  (KNN)-based systems have received the most positive attention because their electromechanical properties can be greatly enhanced using Li and Ta in substitution for K (Na) and Nb, respectively [3–9]. The more general system  $(K,Na,Li)(Nb,Ta,Sb)O_3$ , which not only can be synthesized by the conventional pressureless solid-state sintering method [6–9], but also possesses high piezoelectric  $d_{33}$  values of  $>300$  pC/N in polycrystalline ceramics and  $>400$  pC/N in textured ceramics [10], has been widely accepted as a practicable lead-free replacement for the lead-based systems.

It has been reported that the good electromechanical properties exhibited by the Li- and Ta-modified KNN-based  $[(K,Na,Li)(Nb,Ta)O_3]$  lead-free system are related to the existence of a “morphotropic phase boundary” (MPB) between the orthorhombic (*O*) and tetragonal (*T*) phases [3–10]. In fact, the classical MPB presented in the PZT lead-based system has a narrow composition region, whereby the coexistence of two different phases (or a transitional phase with lower symmetry such as monoclinic) occurs nearly vertical along the temperature scale and the electromechan-

ical properties at a composition in close proximity to that MPB are nearly temperature independent [11]. Here, it turns to two physically interesting and technologically important questions about whether the recently reported *O-T* “MPB” in the  $(K,Na,Li)(Nb,Ta)O_3$  system is a classical MPB similar to that intrinsic in the PZT system, and whether the high electromechanical activity associated with such *O-T* “MPB” has temperature stability as good as that with the classical MPB. In this work, we aim to investigate the temperature dependence of electromechanical properties of a representative  $(K,Na,Li)(Nb,Ta)O_3$  lead-free piezoelectric ceramic system:  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) so as to find out the origin of the phase transition behavior, the nature and validity of the recently reported *O-T* “MPB”, and the temperature stability of the electromechanical activity.

$[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics were synthesized using a conventional solid-state reaction technique [3,10,11]. The compound with  $x=0.03$  was recently reported to appear near the *O-T* “MPB” with good electromechanical properties [6,10]. The starting chemicals were  $K_2CO_3$  (99.5%),  $Na_2CO_3$  (99.9%),  $Li_2CO_3$  (99.9%),  $Nb_2O_5$  (99.9%), and  $Ta_2O_5$  (99.9%). The precursor powder was calcined at  $900^\circ C$  for 2 h in a  $K_2O$ - and  $Na_2O$ -rich atmosphere, and the pressed pellets were sintered at  $1160^\circ C$  for 1 h in air. The typical density of the sintered ceramics was determined to be  $>4300$  kg/m<sup>3</sup> using Archimedes’ method. The temperature dependence of relative permittivity for the unpoled ceramics was evaluated at 1 kHz using a LCR meter (HIOKI 3532) with a temperature chamber. By immersing the ceramics in a silicon oil bath, they were poled with an electric field of 30 kV/cm at  $140^\circ C$  for 30 min before being naturally cooled to room temperature under the electric field. The unipolar electrostrain curves of the poled ceramics

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**Fig. 1.** Temperature dependence of relative permittivity for unpoled  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics at 1 kHz.

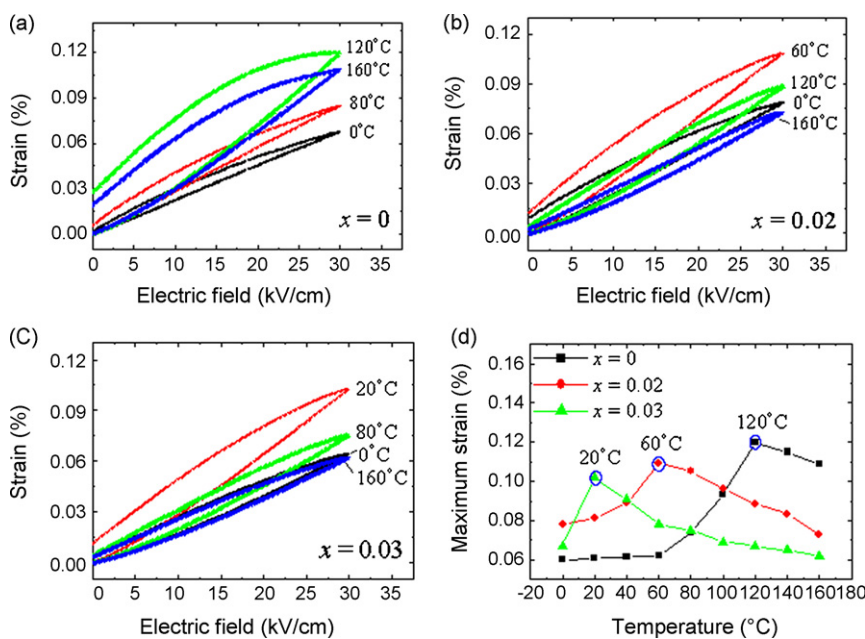
were measured at 5 Hz using a precision ferroelectric test system (Radiant Workstation) and a photonic displacement sensor (MTI 2000) under different temperatures in a temperature-controlled silicon oil bath. The piezoelectric coefficient ( $d_{33}$ ) and electromechanical coupling coefficient ( $k_p$ ) as a function of temperature for the poled ceramics were obtained, respectively, using a  $d_{33}$  meter (IOA-CAS ZJ-3A) and an impedance analyzer (Agilent 4294A), both with a temperature chamber.

**Fig. 1** shows the temperature dependence of relative permittivity for the unpoled  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics at 1 kHz. For the ceramics without Li modification ( $x=0$ ):  $(K_{0.5}Na_{0.5})(Nb_{0.8}Ta_{0.2})O_3$ , two obvious dielectric peaks, which correspond to the phase transitions from  $O$  to  $T$  at  $T_{O-T}=200$  °C and from  $T$  to  $C$  (cubic) at  $T_C=420$  °C in KNN ceramics [3], are found to appear at much lower temperatures of 123 and 280 °C, respectively. By increasing the Li content ( $x=0.02$  and  $0.03$ ),  $T_{O-T}$  shifts downward to the lower temperature side, while  $T_C$  shifts upward to the higher temperature side. It is noted that  $T_{O-T}$  of the  $[(K_{0.5}Na_{0.5})_{0.97}Li_{0.03}](Nb_{0.8}Ta_{0.2})O_3$  ceramics ( $x=0.03$ ) has undergone significant downward shifts to around room temperature of 20 °C in comparison with the  $(K_{0.5}Na_{0.5})(Nb_{0.8}Ta_{0.2})O_3$  ceramics

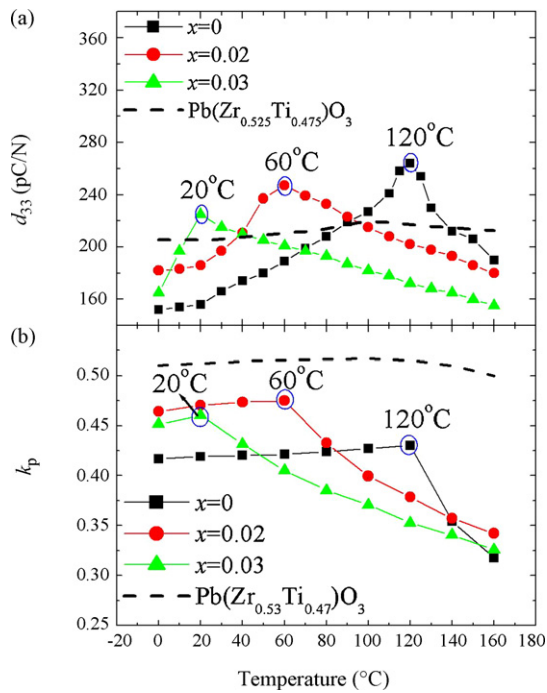
( $x=0$ ) of 123 °C and the  $[(K_{0.5}Na_{0.5})_{0.98}Li_{0.02}](Nb_{0.8}Ta_{0.2})O_3$  ceramics ( $x=0.02$ ) of 62 °C.

**Fig. 2(a)–(c)** illustrates the unipolar electrostrain curves at 5 Hz for the poled  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics at different temperatures of 0–160 °C, whereas **Fig. 2(d)** summarizes the variations of the maximum strain at 30 kV/mm with temperature for the ceramics in **Fig. 2(a)–(c)**. For the  $(K_{0.5}Na_{0.5})(Nb_{0.8}Ta_{0.2})O_3$  ( $x=0$ ) ceramics, the maximum electrostrain of 0.12% at 30 kV/cm is obtained at 120 °C [**Fig. 2(a)** and (d)]. This temperature is referred to  $T_{O-T}=123$  °C in **Fig. 1**, and a slight deviation from this  $T_{O-T}$  leads to a large reduction in the electrostrain values [**Fig. 2(a)** and (d)]. For the  $[(K_{0.5}Na_{0.5})_{0.98}Li_{0.02}](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0.02$ ) ceramics, a reduced maximum electrostrain of 0.11% at 30 kV/cm is observed at their  $T_{O-T}$  of 60 °C [**Fig. 2(b)** and (d)]. Away from this  $T_{O-T}$ , the electrostrain values decrease rapidly. For the  $[(K_{0.5}Na_{0.5})_{0.97}Li_{0.03}](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0.03$ ) ceramics near the recently reported  $O-T$  “MPB”, a further reduced maximum electrostrain of 0.10% at 30 kV/cm is achieved at their  $T_{O-T}$  of 20 °C [**Fig. 2(c)** and (d)]. Similar to the ceramics with  $x=0$  and  $0.02$ , the electrostrain values decrease greatly once the operating temperature is away from this  $T_{O-T}$ . The above observations indicate the presence of a highly temperature-dependent electromechanical behavior at and around  $T_{O-T}$ , besides a common mechanism of temperature impact, in our  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ceramics.

**Fig. 3** plots the piezoelectric coefficient ( $d_{33}$ ) and electromechanical coupling coefficient ( $k_p$ ) as a function of temperature for the poled  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics, together with the corresponding  $d_{33}$  and  $k_p$  values associated with the classical MPBs for the  $Pb(Zr_{0.525}Ti_{0.475})O_3$  [12] and  $Pb(Zr_{0.53}Ti_{0.47})O_3$  [11] ceramics, respectively. As shown in **Fig. 3(a)**,  $d_{33}$  in  $(K_{0.5}Na_{0.5})(Nb_{0.8}Ta_{0.2})O_3$  ( $x=0$ ) ceramics increases gradually from 152 pC/N at 0 °C to 264 pC/N at 120 °C ( $T_{O-T}$ ) and then decreases sharply with increasing temperature to 160 °C. For the  $[(K_{0.5}Na_{0.5})_{0.98}Li_{0.02}](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0.02$ ) ceramics,  $d_{33}$  increases relatively rapidly from 182 pC/N at 0 °C to 247 pC/N at 60 °C ( $T_{O-T}$ ) and then decreases with increasing temperature. For the  $[(K_{0.5}Na_{0.5})_{0.97}Li_{0.03}](Nb_{0.8}Ta_{0.2})O_3$  ( $x=0.03$ ) ceramics near the recently reported  $O-T$  “MPB”,  $d_{33}$  undergoes an increasing and decreasing trend not only similar to the ceramics with



**Fig. 2.** Unipolar electrostrain curves at 5 Hz for poled  $[(K_{0.5}Na_{0.5})_{1-x}Li_x](Nb_{0.8}Ta_{0.2})O_3$  ceramics with (a)  $x=0$ , (b)  $x=0.02$ , and (c)  $x=0.03$  at different temperatures of 0–160 °C. (d) Variations of the maximum strain at 30 kV/mm with temperature for the ceramics in (a)–(c).

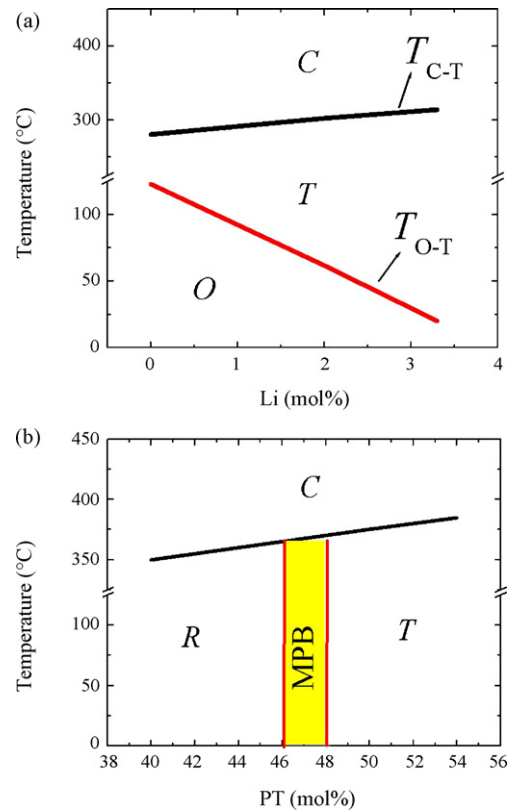


**Fig. 3.** (a) Piezoelectric coefficient ( $d_{33}$ ) and (b) electromechanical coupling coefficient ( $k_p$ ) as a function of temperature for poled  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics, together with the corresponding  $d_{33}$  and  $k_p$  values associated with the classical MPBs for  $\text{Pb}(\text{Zr}_{0.525}\text{Ti}_{0.475})\text{O}_3$  [12] and  $\text{Pb}(\text{Zr}_{0.53}\text{Ti}_{0.47})\text{O}_3$  [11] ceramics, respectively.

$x=0$  and  $0.02$  but also exhibiting the largest  $d_{33}$  of  $225$  pC/N at  $20^{\circ}\text{C}$  ( $T_{O-T}$ ). The temperature dependence of  $d_{33}$  exhibits similar quantitative trends to the temperature dependence of the unipolar electrostrain as shown in Fig. 2(d), suggesting that  $d_{33}$  of our  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ceramics near the recently reported  $O-T$  “MPB” depends strongly on temperature. The  $d_{33}$  result is essentially different from that of the classical MPB for the PZT system where  $d_{33}$  is almost independent of temperature as included in Fig. 3(a) for  $\text{Pb}(\text{Zr}_{0.525}\text{Ti}_{0.475})\text{O}_3$  [12].

Referring to Fig. 3(b),  $k_p$  in the poled  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ceramics with  $x=0$  and  $0.02$  increases slightly with increasing temperature, reaches their maximum values of  $0.43$  and  $0.48$  at their respective  $T_{O-T}$  of  $120$  and  $60^{\circ}\text{C}$ , respectively, and then decreases rapidly with increasing temperature. For the  $[(\text{K}_{0.5}\text{Na}_{0.5})_{0.97}\text{Li}_{0.03}](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ( $x=0.03$ ) ceramics near the recently reported  $O-T$  “MPB”, a similarly large variation in  $k_p$  with temperature is also observed. Although its  $k_p$  can reach a quite high value of  $0.46$  at its  $T_{O-T}$  of  $20^{\circ}\text{C}$ ,  $k_p$  indeed decreases rapidly with increasing temperature above  $T_{O-T}$ . This is also very different from the invariance of  $k_p$  with temperature in the classical MPB for the PZT system as shown in Fig. 3(b) for  $\text{Pb}(\text{Zr}_{0.53}\text{Ti}_{0.47})\text{O}_3$  [11].

The above investigation results clearly show that the electromechanical activity (electrostrain,  $d_{33}$ , and  $k_p$ ) in our  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ( $x=0, 0.02$ , and  $0.03$ ) ceramics increases initially, attains the maximum at their respective  $T_{O-T}$ , and then decreases with increasing temperature. This characteristic trend can basically be preserved irrespective of the compositions, indicating that the maximization of the electromechanical activity at  $T_{O-T}$  is not related to the classical MPB but is attributed to the  $O-T$  ferroelectric phase transition instead. In fact, the ferroelectric-related phase transition is capable of providing a greater contribution to high electromechanical activity due to lattice softening in accordance with the PMN–PT and  $\text{BaTiO}_3$  single crystals [13–16]. Fig. 4 gives the phase diagrams of our  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ceramic system and the PZT



**Fig. 4.** Phase diagrams of (a)  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  and (b)  $\text{Pb}(\text{Zr}_{1-x}\text{Ti}_x)\text{O}_3$  (PZT) ceramic systems constructed from the phase transition temperatures  $T_{O-T}$  and  $T_C$  of Fig. 1 and the relevant data of Ref. [12], respectively. C = cubic, T = tetragonal, O = orthorhombic, and R = rhombohedral.

ceramic system, which are constructed from the phase transition temperatures  $T_{O-T}$  and  $T_C$  of Fig. 1 and the relevant data of Ref. [12], respectively. In Fig. 4(a), an increase in Li modification leads to a shift in  $T_{O-T}$  to a lower temperature, but it is unable to constitute a classical MPB like that in the PZT system. It should be noted that the classical MPB for the PZT system represents a narrow composition region with two different phases coexisting nearly vertical along the temperature scale, as illustrated in Fig. 4(b). Therefore, the recently reported  $O-T$  “MPB” in the  $(\text{K},\text{Na},\text{Li})(\text{Nb},\text{Ta})\text{O}_3$  system is definitely not a classical MPB engaged in the PZT system; it should be an  $O-T$  ferroelectric phase transition region at ambient temperature. More recent studies on X-ray diffraction data and Raman scattering spectra of the  $(\text{K},\text{Na},\text{Li})(\text{Nb},\text{Ta})\text{O}_3$  system support our findings from microscopic perspectives [17,18]. Away from  $T_{O-T}$ , the phase transition and related domain reorientation in the ceramics result in a rapid decrease in their high electromechanical activity, giving rise to strong temperature dependence in general.

In conclusion, temperature dependence of electromechanical properties of  $[(\text{K}_{0.5}\text{Na}_{0.5})_{1-x}\text{Li}_x](\text{Nb}_{0.8}\text{Ta}_{0.2})\text{O}_3$  ( $x=0, 0.02$ , and  $0.03$ ) lead-free ceramics has been investigated. It has been found that the electromechanical activity of the ceramics increases rapidly to the maximum at  $T_{O-T}$  as a result of the existence of an  $O-T$  phase transition region at ambient temperature in opposition to a classical MPB region as commonly seen in the PZT system. Moreover, Li modification effectively shifts  $T_{O-T}$  to lower temperatures, but ineffectively constitutes a classical MPB. The  $O-T$  “MPB” as per the recent reports is only an  $O-T$  ferroelectric phase transition region at ambient temperature. The strong temperature dependence of electromechanical activity in the ceramics may limit their application viability.

## Acknowledgments

This work was supported by the Innovation and Technology Fund of the HKSAR Government under Grant No. GHP/003/06.

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